

# Tight reference frame-independent quantum teleportation

Dominic Verdon & Jamie Vicary

Department of Computer Science  
University of Oxford

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- Depends on the action of the group of RF transformations.  
The group must be **finite**.
- Constructions of reference frame-independent (RFI) teleportation protocols.

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- Work has been done on imperfect protocols in the infinite case. [Marzolino and Buchleitner, 2015]
- We deal with the case of a finite group of RF transformations.

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  - Imperfect protocols as limits of perfect finite-group schemes?
- CQM as a way to work with RFs in quantum information.

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  - Reference frames
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- $H$  carries a unitary representation  $\pi : G \rightarrow \text{End}(H)$  of  $G$ .
- When RF configuration transforms by  $g^{-1} \in G$ :

	Old frame	New frame
State	$ \psi\rangle$	$\pi(g) \psi\rangle$
Operations	$L \in \text{End}(H)$	$\pi(g)L\pi(g)^\dagger$

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  - 1  $U_i$  are all unitary.
  - 2  $U_i$  are orthonormal under the Hilbert-Schmidt inner product.
- When Hilbert spaces have minimal dimension, teleportation protocols correspond to UEBs:

Shared entangled state	$\sum_i  i\rangle \otimes  i\rangle$
Alice's measurement basis	$ \phi_x\rangle := \sum_i  i\rangle \otimes U_x  i\rangle$
Bob's unitary correction	$C_x := U_x^T$

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## Example: II

- Take  $G = \mathbb{Z}^2$ , where the nontrivial element  $a \in G$  acts as

$$\pi(a) = \begin{pmatrix} \sqrt{3}/2 & 1/2 \\ 1/2 & -\sqrt{3}/2 \end{pmatrix}.$$



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- They agree to use the following UEB:

$$U_0 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} \quad U_2 = \frac{1}{4} \begin{pmatrix} -\sqrt{2} - \sqrt{6} & -\sqrt{2} + \sqrt{6} \\ -\sqrt{2} + \sqrt{6} & \sqrt{2} + \sqrt{6} \end{pmatrix}$$

$$U_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \quad U_3 = \frac{1}{4} \begin{pmatrix} \sqrt{2} - \sqrt{6} & -\sqrt{2} - \sqrt{6} \\ -\sqrt{2} - \sqrt{6} & -\sqrt{2} + \sqrt{6} \end{pmatrix}$$

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  - Alice's perspective: Bob's correction was *not*  $U_i^T$  but rather  $\pi(a)^\dagger U_i^T \pi(a)$ .
  - Bob's perspective: the measurement result  $i$  Alice communicated did not correspond to the state  $(\mathbb{1} \otimes \pi(a)) |\phi_i\rangle$  she measured.

# The failure of speakable communication: a theorem

## Theorem (VV)

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- Teleportation is possible if and only if  $\pi(g)^\dagger U_i^T \pi(g) = U_i^T$  for all  $i$  and  $g \in G$ .



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- Therefore all irreducible factors of  $H$  are identical and 1D.



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- Suppose Bob's lab is aligned upside-down wrt Alice's. 0, 1, 2 or 3 will be received as 1, 0, 3 or 2 respectively.

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- But the UEB is such that the communication error cancels with his correction error:

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- So teleportation is successful regardless of RF alignment!



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  - ② The classical channel carried a  $G$ -action, so we could encode the measurement results to carry the **inverse** permutation to the UEB.
- If we can find a suitable classical channel and a  $G$ -equivariant UEB, we can perform RFI teleportation.

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### Theorem (VV)

*This procedure is RF-independent exactly when the UEB is  $G$ -equivariant.*



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- We provide methods for constructing them when they do.

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- $G$ -equivariant orthonormal bases can be easily classified!

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Two coset spaces are isomorphic  
 $\iff$  they correspond to conjugate subgroups.
- $\mathcal{M}$  is additive ( $\sqcup \rightarrow \oplus$ ).
- So in order to classify all objects in  $\text{Im}(\mathcal{M})$ , it is sufficient to find the images of the coset spaces  $G/H$  under  $\mathcal{M}$  - the **basic permutation representations**.

# A representation on which no $G$ -equivariant UEBs exist

## Theorem (VV)

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- This does not work for all irreps! (2D irrep of  $D_8$ .)

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## Theorem (VV)

Let  $|v_i\rangle$  be the G-equivariant orthonormal basis, and H be the Hadamard matrix. Then the G-equivariant UEB is:

$$(U_H)_{ij} = \frac{1}{N} H \circ \text{diag}(H, j)^\dagger \circ H^\dagger \circ \text{diag}(H^T, i) \quad (1)$$



# A sufficient condition for $\dim < 5$

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- Any  $M < C(S_d)$  has identical entries on the diagonal and identical entries everywhere else.

# A sufficient condition for $\dim < 5$ (cont.)

## Proof (Cont.)

Unitarity of such a matrix is equivalent to

$$|b|^2 = \frac{1 - |a|^2}{d - 1} \quad (2)$$

$$\Re(a^*b) = \frac{2 - d}{2} |b|^2. \quad (3)$$

These equations can be satisfied for  $|a|, |b| = \frac{1}{\sqrt{d}}$  iff  $d < 5$ .  $\square$

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- In a dagger-compact category, a *quantum teleportation procedure* is a classical structure on  $A \otimes A^*$  satisfying:

$$\begin{array}{c} \uparrow \\ \uparrow \\ \text{comultiplication} \\ \downarrow \\ \downarrow \end{array} \begin{array}{c} \downarrow \\ \downarrow \\ \uparrow \\ \uparrow \end{array} = c \cdot \begin{array}{c} \uparrow \\ \uparrow \\ \text{unit} \\ \downarrow \\ \downarrow \end{array} \begin{array}{c} \downarrow \\ \downarrow \\ \uparrow \\ \uparrow \end{array} \quad (4)$$

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- In  $\mathbf{Rep}(G)$  these are exactly the  $G$ -equivariant unitary error bases.
- All categorical constructions of UEBs carry over to  $\mathbf{Rep}(G)$ .
- Unclear how to generalise non-categorical constructions to the  $G$ -equivariant setting.

# Future work

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- Teleportation of anyons.

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